# Pressure-Volume-Temperature Behavior of Nylon 610 

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## Synopsis

Equilibrium pressure-volume-temperature behavior in both the solid and molten regions was determined for nylon 610. Data were measured with a compressibility device capable of obtaining precise and accurate data. Residual curve treatment showed that the data were true equilibrium data. A volume extrapolated to 2419 atm . at room temperature from the present data compared favorably to the sole literature value reported by Bridgman. The data of this work showed the existence of what appears to be a sec-ond-order transition point for nylon 610 . This point ranged from $140^{\circ} \mathrm{C}$. at 232 atm . to about $170^{\circ} \mathrm{C}$. at 1855 atm . The Spencer-Gilmore equation was fitted to the data of this study.

Investigations on the effect of pressure on the physical and chemical behavior of polymers are subjects of interest to both polymer engineers and scientists. An area of particular interest is that involving the pressure-volume-temperature relations of polymers.

A number of investigators, Bridgman, ${ }^{1}$ Parks and Richards, ${ }^{2}$ Spencer and Gilmore, ${ }^{3,4}$ Weir, ${ }^{5-7}$ Matsuoka, ${ }^{8-10}$ Heydemann and Guicking, ${ }^{11}$ Hellwege et al. ${ }^{12}$ and Foster et al. ${ }^{13,14}$ have measured pressure-volume-temperature data. Many of these studies have, however, been limited with respect to temperature, pressure, or types of polymers studied. In particular nylon 610 is one polymer that has been neglected. The only available literature data is one value of compressibility. ${ }^{1}$

The present work was undertaken to study as completely as possible both the solid and melt compressibility behavior of nylon 610.
Figure 1 is a schematic of the experimental apparatus. Figure 2 depicts the compressibility chamber and piston. Polymer samples were placed in the chamber which was evacuated after it had been immersed in the oil bath thermostat. The device was then thoroughly seasoned. A careful annealing procedure was followed for solid phase measurements. Readings were then taken by holding temperature constant, changing volume, and then balancing the pressure for a long enough period of time to assure the proper equilibrium pressure-volume-temperature combination. A detailed description of apparatus and procedure is given elsewhere. ${ }^{15}$

Temperatures were controlled to $\pm 0.10^{\circ} \mathrm{C}$. Volume accuracy was $0.25 \%$ or better, while the accuracy of pressure measurements was within $1 \%$ or less.

[^0]Data were measured for both the solid and molten regions from pressures of $1-1855 \mathrm{~atm}$. and temperatures of $23-260^{\circ} \mathrm{C}$. These raw data were smoothed by using volume and pressure residual curves. The technique of smoothing with residuals involves first calculating either a volume or


Fig. 1. Schematic of apparatus.
pressure from an equation of state (at a selected temperature and pressure or a selected volume and pressure). In this work the Spencer and Gilmore ${ }^{3,4}$ equation of state was used:

$$
\begin{equation*}
(P+\pi)(V-W)=(R / M) T \tag{1}
\end{equation*}
$$

where $P$ is pressure, $V$ is specific volume, $R$ is the gas constant, $T$ is absolute temperature, $M$ is the molecular weight of an interaction unit, and


Fig. 2. Compressibility chamber.


Fig. 3. Volume residual vs. pressure.


Fig. 4. Pressure vs. volume at constant temperature.
$\pi$ and $W$ are constants. Next the pressure or volume residual is calculated.

$$
\begin{align*}
& r_{p}=P_{\text {calculated }}-P_{\text {experimental }}  \tag{2}\\
& r_{v}=V_{\text {calculated }}-V_{\text {experimental }} \tag{3}
\end{align*}
$$

These residual values can then be plotted as isothermal curves versus either volume or pressure. Figure 3 is such a plot. Smoothed data points can then be obtained by using the residual together with calculated values:

$$
\begin{align*}
P_{\text {experimental }} & =P_{\text {calculated }}-r_{p}  \tag{4}\\
V_{\text {experimental }} & =V_{\text {calculated }}-r_{v} \tag{5}
\end{align*}
$$

The smoothed data are presented in Figures 4 and 5 and Table I.
TABLE I
Smoothed Values of the Specific Volume of Nylon 610 for Temperatures of $25-260^{\circ} \mathrm{C}$. and Pressures of $1-1855 \mathrm{~atm}$.

| Temperature, ${ }^{\circ} \mathrm{C}$. | Specific volume, cm. ${ }^{3} / \mathrm{g} . \times 10^{4}$ |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | 1 atm . | 79 | 158 | 232 | 316 | 474 | 618 | 1242 | 1855 |
| 25 | 9226 | 9198 | 9182 | 9161 | 9140 | 9097 | 9040 | 8898 | 8881 |
| 66 | 9296 | 9266 | 9231 | 9206 | 9178 | 9131 | 9084 | 8963 | 8948 |
| 93 | 9324 | 9292 | 9260 | 9233 | 9197 | 9151 | 9113 | 9015 | 9000 |
| 121 | 9350 | 9333 | 9307 | 9276 | 9255 | 9212 | 9170 | 9048 | 9032 |
| 149 | 9364 | 9362 | 9361 | 9359 | 9355 | 9347 | 9331 | 9167 | 9094 |
| 177 | 9382 | 9379 | 9377 | 9375 | 9373 | 9365 | 9356 | 9300 | 9200 |
| 199 | 9391 | 9389 | 9386 | 9384 | 9381 | 9374 | 9366 | 9306 | 9205 |
| 204 | 9401 | 9396 | 9394 | 9392 | 9389 | 9381 | 9371 | 9312 | 9220 |
| 21.0 | 9420 | 9417 | 9412 | 9409 | 9405 | 9396 | 9384 | 9320 | 9225 |
| 216 | 9459 | 9456 | 9451 | 9447 | 9441 | 9424 | 9408 | 9327 | 9227 |
| 221 | 9582 | 9579 | 9574 | 9567 | 9557 | 9532 | 9506 | 9373 | 9268 |
| 227 | 9778 | 9753 | 9723 | 9695 | 9663 | 9578 | 9528 | 9391 | 9282 |
| 232 | 9865 | 9832 | 9774 | 9727 | 9679 | 9587 | 9530 | - | - |
| 260 | 9874 | - | - | - | - | - | - | - | - |



Fig. 5. Volume vs. temperature at constant pressure.

Residual curves can also be used to check the consistency of experimental data as well as the attainment of equilibrium during the experiments. ${ }^{16,17}$ Smooth curves, as in Figure 3, indicate that data are consistent and that equilibrium was attained.

The only other compressibility datum for nylon 610 was a single point reported by Bridgman ${ }^{1}$ at room temperature and a pressure of 2419 atm . A direct comparison of the data of this work and the point of Bridgman ${ }^{1}$ was not possible, since the highest pressure studied in the present investigation was 1855 atm . It was found, however, that a volume extrapolated
to 2419 atm . at $25^{\circ} \mathrm{C}$. from the data of the present work compared quite favorably to the value reported at that pressure by Bridgman. ${ }^{1}$

The curves of Figure 5 show two distinct discontinuities. One of these, which appears in the range of $200-220^{\circ} \mathrm{C}$., represents the melting of the sample. There is, however, another which appears from about $140^{\circ} \mathrm{C}$. at 232 atm . to about $170^{\circ} \mathrm{C}$. at 1855 atm . This latter discontinuity is not as readily explained.

The most obvious explanation for the discontinuity is that it might be due to experimental error. This, however, did not seem to be the case since a number of independent rechecks yielded essentially the same behavior.

Another possibility is that the discontinuity might be the glass transition point. Nielsen, ${ }^{18}$ however, reports the $T_{g}$ for nylon 610 as $40^{\circ} \mathrm{C}$., which discounts this explanation. It appears more likely that the discontinuity represents a second-order transition in the polymer.

Shifts of both glass transition points and second-order transitions with increasing pressure have been observed previously by other investigators. ${ }^{11}$ These shifts ranged from 0.013 to $0.018^{\circ} \mathrm{C} . / \mathrm{atm}$. pressure. In the present work the shift was about $0.016^{\circ} \mathrm{C} . / \mathrm{atm}$. which was approximately the same as those observed earlier. ${ }^{11}$

The Spencer-Gilmore equation was fitted to the data of the present work in the solid region. Two sets of constants were determined-one below the second-order transition point and the other between that point and the melting region. These constants are given in Table II.

TABLE II

| $R^{\prime}, \mathrm{cm}^{3}-\mathrm{atm} . / \mathrm{g} .^{\circ} \mathrm{K}^{a}$ | $\pi, \mathrm{~atm}$. | $W, \mathrm{~cm}^{3} / \mathrm{g}$. | Temperature <br> range, ${ }^{\circ} \mathrm{C}$. |
| :---: | ---: | :---: | :---: |
| 1.85 | 1,351 | 0.8603 | $25-121$ |
| 0.738 | 10,768 | 0.9064 | $177-221$ |

${ }^{\text {a }} R^{\prime}=R / M$.
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## Résumé

Le comportement pression-volume-température à l'équilibre à la fois dans les régions du solide et liquide a été déterminé pour du nylon 610. Les résultats ont été obtenus grâce à un appareillage de mesures de compressibilité permettant l'obtention de données précises. Le traitement de la courbe montrait que ces résultats correspondaient à un véritable équilibre. Le volume-extrapolé à 2.409 atmosphères à température de chambre au départ des résultats actuels peuvent être comparés favorablement à la seule valeur indiquée dans la littérature et qui est rapportée par Bridgman. Les résultats de ce travail montraient l'existence de ce qui semblait être un point de transition de second ordre pour le nylon 610. Ce point se situe entre $140^{\circ}$ sous 232 atmosphères jusqu'à environ $170^{\circ} \mathrm{C}$ sous 1.855 atmosphères. L'équation de Spencer-Gilmore était en accord avec les résultats de cette étude.

## Zusammenfassung

Das Gleichgewichtsvolumen von Nylon 610 in Abhängigkeit von Druck und Temperatur wurde in geschmolzener und fester Phase bestimmt. Die Messung erfolgte mit einem. Kompressibilitätsapparat, der Präzisionsdaten lieferte. Das aus den erhaltenen Daten bei Raumtemperatur auf 2419 Atmosphären extrapolierte Volumen stimmte gut mit dem einzigen, in der Literatur vorhandenen, von Bridgman bestimmten Wert überein. Die Ergebnisse der vorliegenden Arbeit lassen anscheinend das Vorhandensein eins Umwandlungspunkts zweiter Ordnung bei Nylon 610 erkennen. Dieser Punkt verschob sich von $140^{\circ} \mathrm{C}$ bei 232 Atmosphären zu etwa $170^{\circ} \mathrm{C}$ bei 1855 Atmosphären. Die Ergebnisse konnten mit der Spencer-Gilmore-Gleichung wiedergegeben werden.

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